

# Nonlinear elliptic involving the $p$ -Laplacian problems and stability of positive solutions

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**Abstract.** This paper deal with the stability of positive stationary solutions of boundary value problem of the form

$$\begin{cases} -\Delta_p u^m(x) = \lambda a(x)f(u(x)), & x \in \Omega \\ u(x) = 0, & x \in \partial\Omega \end{cases}$$

where  $\Omega \subset R^N (N \geq 1)$  is a bounded domain with a smooth boundary  $\partial\Omega$ ,  $\Delta_s z = \text{div}(|\nabla z|^{s-2} \nabla z)$  is the  $p$ -Laplacian operator;  $s > 2$ ,  $\lambda > 0$  is a parameter,  $m > 1$ ,  $a(x)$  is a weight function that satisfies  $a(x) \in C(\bar{\Omega})$  for all  $x \in \Omega$ . Here  $f : [0, +\infty) \rightarrow R^-(R^+)$  is  $C^1$  function. We establish if  $\frac{f(u)}{u^{m(p-1)}}$  be strictly decreasing (increasing) then every positive stationary solution is stable.

**Keywords:** nonlinear elliptic,  $p$ -Laplacian problems, stability of positive solutions

## 1 Section heading

We will Consider the following  $p$ -Laplacian equation:

$$-\Delta_p u^m(x) = \lambda a(x)f(u(x)), \quad x \in \Omega \quad (1)$$

$$u(x) = 0, \quad x \in \partial\Omega \quad (2)$$

where  $\Delta_s z = \text{div}(|\nabla z|^{s-2} \nabla z)$  denotes the  $p$ -Laplacian operator;  $s > 2$  and  $\lambda$  is a positive parameter,  $m > 1$  is a constant,  $\Omega \subset R^N (N \geq 1)$  is a bounded region with smooth boundary  $\partial\Omega$ , weight  $a(x)$  satisfies  $a(x) \in C(\bar{\Omega})$  and we assume that either  $a(x) > 0$  or  $a(x) < 0$  for all  $x \in \Omega$ .

Existence of solutions in this type of models (the case system,  $p = 2$ ) have been studied in [3]. A large number of works have been made studying stability of solutions in the case when  $m = 1$  (see [1, 2, 4, 5]). In [1], authors have shown that every non-negative stationary solution of boundary value problem

$$\begin{cases} -\Delta_p u(x) = \lambda g(x)f(u(x)), & x \in \Omega \\ Bu(x) = 0, & x \in \partial\Omega \end{cases}$$

is linearly unstable (stable). But in this article, we ponder the case when  $m > 1$  and establish stability of positive stationary solutions of (1) and (2), under certain conditions.

It is well known that if  $u$  be any positive solution of (1) and (2), then the linearized equation about  $u$  consists of

$$-m(p-1)\text{div}(|\nabla u^m|^{p-2} \nabla(u^{m-1}\phi)) - \lambda a(x)f'(u)\phi = \mu\phi, \quad x \in \Omega \quad (3)$$

$$\phi(x) = 0, \quad x \in \partial\Omega \quad (4)$$

where equation (3) obtained from the formal derivative of the operator  $\Delta_p u^m$  with respect to  $u$ .

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**Definition 1.** We say a positive solution  $u$  is stable if all eigenvalues of (3) is strictly positive, which can be inferred if the principal eigenvalue  $\mu_1$  (i.e., eigenvalue corresponding to positive eigenfunction) be positive. Otherwise we say  $u$  is unstable.

## 2 Main results

In this section, we prove the following Theorem by determining the sign of the principal eigenvalue,  $\mu_1$ , of the linearized equation of (1) about a positive solution  $\bar{u}$ . Stability of solutions follows from Theorem (4.2) of [6] that if  $\mu_1 > 0$ , then  $\bar{u}$  is stable, otherwise  $\bar{u}$  is unstable.

**Theorem 1.** Let  $f(u): [0, \infty) \rightarrow R^-(R^+)$ ,  $a(x) > 0(a(x) < 0)$ , and  $\frac{f(u)}{u^{m(p-1)}}$  be strictly decreasing (increasing), then every positive stationary solution of boundary value problem (1) and (2) is linearly stable.

*Proof.* Let  $\bar{u}$  be any positive stationary solution of (1) and (2), then the linearized equation about  $\bar{u}$  consists of

$$-m(p-1)\operatorname{div}(|\nabla\bar{u}^m|^{p-2}\nabla(\bar{u}^{m-1}\phi)) - \lambda a(x)f'(\bar{u})\phi = \mu\phi, \quad x \in \Omega \quad (5)$$

$$\phi(x) = 0, \quad x \in \partial\Omega \quad (6)$$

presume that  $\mu_1$  is the principal eigenvalue and also  $\psi(x)(\psi \geq 0)$  is a corresponding eigenfunction. Now multiplying (5) by  $-\bar{u}$ , (1) by  $m(p-1)\psi$ , add integrating by parts over  $\Omega$  and add the two resulting expressions to get

$$\begin{aligned} -\mu_1 \int_{\Omega} \psi(x)\bar{u}(x)dx &= m(p-1) \int_{\Omega} [\bar{u}\operatorname{div}(|\nabla\bar{u}^m|^{p-2}\nabla(\bar{u}^{m-1}\psi)) - \psi\operatorname{div}(|\nabla\bar{u}^m|^{p-2}\nabla\bar{u}^m)]dx \\ &\quad + \lambda \int_{\Omega} \psi(x)a(x)[\bar{u}f'(\bar{u}) - m(p-1)f(\bar{u})]dx. \end{aligned} \quad (7)$$

But by Green identity and Dirichlet boundary condition,

$$\begin{aligned} \int_{\Omega} \bar{u}\operatorname{div}(|\nabla\bar{u}^m|^{p-2}\nabla(\bar{u}^{m-1}\psi))dx &= \int_{\Omega} \bar{u}\nabla(|\nabla\bar{u}^m|^{p-2})\nabla(\bar{u}^{m-1}\psi)dx + \int_{\Omega} \bar{u}|\nabla\bar{u}^m|^{p-2}\Delta(\bar{u}^{m-1}\psi)dx \\ &= - \int_{\Omega} |\nabla\bar{u}^m|^{p-2}\nabla\bar{u}\nabla(\bar{u}^{m-1}\psi)dx, \end{aligned} \quad (8)$$

also similarly,

$$\int_{\Omega} \psi\operatorname{div}(|\nabla\bar{u}^m|^{p-2}\nabla(\bar{u}^m))dx = - \int_{\Omega} |\nabla\bar{u}^m|^{p-2}\nabla\psi\nabla(\bar{u}^m)dx, \quad (9)$$

and hence

$$\begin{aligned} &\int_{\Omega} [\bar{u}\operatorname{div}(|\nabla\bar{u}^m|^{p-2}\nabla(\bar{u}^{m-1}\psi)) - \psi\operatorname{div}(|\nabla\bar{u}^m|^{p-2}\nabla\bar{u}^m)]dx \\ &= \int_{\Omega} |\nabla\bar{u}^m|^{p-2}(\nabla\psi\nabla\bar{u}^m - \nabla\bar{u}\nabla(\bar{u}^{m-1}\psi))dx \\ &= \left(\frac{m-1}{m}\right) \int_{\Omega} |\nabla\bar{u}^m|^{p-2}\nabla\psi\nabla\bar{u}^m dx \\ &\quad - (m-1) \int_{\Omega} |\nabla\bar{u}^m|^{p-2}\bar{u}^{m-2}(\nabla\bar{u})^2\psi dx, \end{aligned} \quad (10)$$

on the other hand by Divergence theorem and since  $\psi$  is eigenfunction (thus  $\psi = 0$  on  $\partial\Omega$ ), we have

$$\int_{\Omega} |\nabla\bar{u}^m|^{p-2}\nabla\bar{u}^m\nabla\psi dx = \int_{\Omega} \operatorname{div}(|\nabla\bar{u}^m|^{p-2}\nabla\bar{u}^m\psi)dx - \int_{\Omega} \operatorname{div}(|\nabla\bar{u}^m|^{p-2}\nabla\bar{u}^m)\psi dx$$

$$\begin{aligned}
 &= \int_{\partial\Omega} |\nabla \bar{u}^m|^{p-2} \nabla \bar{u}^m \psi \cdot n ds - \int_{\Omega} \operatorname{div}(|\nabla \bar{u}^m|^{p-2} \nabla \bar{u}^m) \psi dx \\
 &= \lambda \int_{\Omega} \psi(x) a(x) f(\bar{u}) dx,
 \end{aligned}
 \tag{11}$$

therefore by replacing (11) in (10), (7) becomes

$$\begin{aligned}
 -\mu_1 \int_{\Omega} \psi(x) \bar{u}(x) dx &= (p-1)(m-1) \int_{\Omega} \psi [\lambda a(x) f(\bar{u}) - m |\nabla \bar{u}^m|^{p-2} \bar{u}^{m-2} (\nabla \bar{u})^2] dx \\
 &+ \lambda \int_{\Omega} \psi(x) a(x) [\bar{u} f'(\bar{u}) - m(p-1) f(\bar{u})] dx.
 \end{aligned}
 \tag{12}$$

Now observe that since  $\frac{f(\bar{u})}{\bar{u}^{m(p-1)}}$  is strictly decreasing (increasing),  $\bar{u} > 0$  and  $a(x) > 0 (< 0)$ , we get

$$a(x) [\bar{u} f'(\bar{u}) - m(p-1) f(\bar{u})] < 0,
 \tag{13}$$

and because of  $f(\bar{u}) < 0 (> 0)$ ,  $a(x) > 0 (< 0)$  and  $\bar{u} > 0$ , we have

$$\lambda a(x) f(\bar{u}) - m |\nabla \bar{u}^m|^{p-2} \bar{u}^{m-2} (\nabla \bar{u})^2 < 0,
 \tag{14}$$

thus using (13) and (14) in (12) and knowing this fact that  $\psi > 0$  in  $\Omega$ , we obtain

$$-\mu_1 \int_{\Omega} \psi(x) \bar{u}(x) dx < 0.
 \tag{15}$$

Hence  $\mu_1 > 0$  and it show that every positive solution is stable(see[6]).

Next Theorem can be generalization of Theorem 1.

**Theorem 2.** Suppose  $f(x, u): \Omega \times [0, \infty) \rightarrow R^-$  be a continuous function and  $\frac{f(x,u)}{u^{m(p-1)}}$  be strictly decreasing for  $u > 0$  at each fixed  $x \in \Omega$ , then every positive stationary solution of boundary value problem

$$-\Delta_p u^m(x) = \lambda f(x, u), \quad x \in \Omega
 \tag{16}$$

$$u(x) = 0, \quad x \in \partial\Omega
 \tag{17}$$

is linearly stable.

*Proof.* The proof proceeds in the same way as for Theorem 1, i.e; let  $\bar{u}$  be any positive stationary solution of (16) and (17). Instead of (12), we obtain

$$\begin{aligned}
 -\mu_1 \int_{\Omega} \psi(x) \bar{u}(x) dx &= (p-1)(m-1) \int_{\Omega} \psi [\lambda f(x, \bar{u}) - m |\nabla \bar{u}^m|^{p-2} \bar{u}^{m-2} (\nabla \bar{u})^2] dx \\
 &+ \lambda \int_{\Omega} \psi(x) [\bar{u} f'(x, \bar{u}) - m(p-1) f(x, \bar{u})] dx.
 \end{aligned}$$

and result follows via the appropriate analogue of Theorem 1.

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